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# Synthesis, structural and spectroscopic properties of orthorhombic compounds BaLnCuS<sub>3</sub> (Ln = Pr, Sm)

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#### Abstract

Ternary sulfides BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub> are first synthesized by the sulphidation reaction of a mixture of related oxides and metal Cu in a flow of (CS<sub>2</sub>, H<sub>2</sub>S) at 1170 K. The crystal structures of BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub> are obtained by Rietveld method. BaPrCuS<sub>3</sub> crystallizes in space group *Pnma* with unit cell parameters a = 10.56074(6), b = 4.11305(2) and c = 13.42845(7) Å, V = 583.289 (5) Å<sup>3</sup>, Z = 2. BaSmCuS<sub>3</sub> crystallizes in space group *Cmcm* with unit cell parameters a = 4.07269(4), b = 13.4499(1) and c = 10.3704(1) Å, V = 568.06 (1) Å<sup>3</sup>, Z = 2. The structural model is proposed for the *Cmcm*→*Pnma* transition in compounds ABCX<sub>3</sub> (X = S, Se) for the sequence Sm-Pm-Nd-Pr. The dimensionless tolerance factor  $t = IR(A) \times IR(C)/IR(B)^2$  is suggested to control the boundary between the *Cmcm* and *Pnma* structures. The micromorphological, thermal and spectroscopic properties are evaluated for BaPrCuS<sub>3</sub>. The compound melts incongruently at T<sub>melt</sub> = 1580.9 K. In BaPrCuS<sub>3</sub>, the band gap is estimated to be 2.1 eV. The vibrational parameters of BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub> are comparatively observed by Raman spectroscopy.

Keywords: Complex sulfides, Crystal structure, SEM, Raman.

# 1. Introduction

Complex chalcogenide compounds have interesting chemical and physical properties and the materials are widely used in modern optical and electronic technologies [1-10]. In recent years, complex Cu<sup>+</sup>-containing sulfides and selenides were considered as promising thin film absorber materials in solar cell structures [11-16]. This activated general interest to crystal chemistry and the properties of Cu<sup>+</sup> chalcogenides, including a search for new materials. In particular, the compounds with generalized composition ABCX<sub>3</sub>, where A –alkaline-earth metal, B – d- or f-element, C – other d-element, X – chalcogenide, can be considered [1]. Earlier, besides crystal structures, the semiconductor, magnetic, optical and thermodynamic properties were reported on for the selected compounds from this family [1,17-25]. Besides, several compounds  $Eu^{2+}BC'S_3$  have been recently discovered [7,26-28]. It is known that compounds ABCX<sub>3</sub> can crystallize in seven structural types: KZrCuS<sub>3</sub> (*Cmcm*), Eu<sub>2</sub>CuS<sub>3</sub> (*Pnma*), Ba<sub>2</sub>MnS<sub>3</sub> (*Pnma*), BaCuLaS<sub>3</sub> (*Pnma*), BaAgErS<sub>3</sub> (*C2/m*), NaCuTiS<sub>3</sub> (*Pnma*), TlCuTiTe<sub>3</sub> (*P2*<sub>1</sub>/*m*) [1]. Among these crystals, the ALnCuQ<sub>3</sub> compounds are particularly interesting because Ln elements possess rich crystal chemistry and specific spectroscopic properties.

The known ALnCuS<sub>3</sub> compounds crystallize in Eu<sub>2</sub>CuS<sub>3</sub> (*Pnma*) or KZrCuS<sub>3</sub> (*Cmcm*) structural types and the rules governing the relation between the element composition and structure type was not defined for this ternary sulfide family [1]. The crystal structures related to these structural types are very similar. The structures are formed by layers of CuX4 tetrahedra and LnX6 octahedra [29,30]. The present study is aimed at the synthesis and evaluation of compounds BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub> as the representative members of ALnCuS<sub>3</sub> sulfides. The compounds were not observed earlier. As to other BaLnCuS<sub>3</sub> sulfides, several compounds were synthesized in the past for Ln= La, Ce, Nd, Er, Sc, Y, and the structures were defined for BaLaCuS<sub>3</sub> (*Pnma*) and BaErCuS<sub>3</sub> (*Cmcm*) [17,18]. Also, the space groups *Pnma* and *Cmcm* were adopted for Ln = Ce, Nd and Ln = Sc, Y, Gd, respectively. Thus, the phase transition appears on the rare-earth element substitution in BaLnCuS<sub>3</sub> structures. The formation of BaCuLnS<sub>3</sub> (Ln = Pr, Sm) is assumed in the quasithreefold BaS-Ln<sub>2</sub>S<sub>3</sub>-Cu<sub>2</sub>S systems. It should be mentioned that the BaS-Ln<sub>2</sub>S<sub>3</sub>-Cu<sub>2</sub>S phase diagram has not been defined yet for Ln = Pr, Sm. The known analogous systems (Ba,Sr,Eu<sup>2+</sup>)S-Ln<sub>2</sub>S<sub>3</sub>-Cu<sub>2</sub>S have the composite equilibriums including peritectic/eutectic interactions and extensive ranges of solid solutions [22,23,28,31,32] and, respectively, the synthesis of BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub> seems to be not trivial. In the present experiment, the sulfides are prepared by the sulphidation method and final

powder products are observed by conventional experimental techniques to see their structural, morphological and vibrational characteristics.

#### 2. Experimental section

**Syntheses**. BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub> compounds were prepared in powder forms by the sulphidation of oxide mixtures received after the decomposition of metal nitrate solutions. In the synthesis, for better comparison, the key steps were the same for both compounds of barium, rareearth element and copper.

The high purity starting reagents were used: Cu (99,9%, SZB Tsvetmet, Russia), BaCO<sub>3</sub> (99,99%, ultrapure, KAI YONG, China),  $Pr_6O_{11}$  (99,99%, ultrapure, TDM-96 Ltd. Russia), Sm<sub>2</sub>O<sub>3</sub> (99,99%, ultrapure, TDM-96 Ltd. Russia) and concentrated nitric acid solution (C(HNO<sub>3</sub>) = 14.6 mol/l, ultrapure, Vekton Ltd., Russia). Ammonium rodanide NH<sub>4</sub>SCN (98%, Vekton Ltd., Russia). Weighing them was carried out on the assay balances of Mettler Toledo with the accuracy of 0.1 mg. Before weighing, a copper plate of ~1 mm in thickness was etched in the HClO<sub>3</sub> solution to remove the surface oxide and was washed out in distilled water. Then, the Cu plate was cut into segments of ~1-5 mm<sup>2</sup> in area. All starting metal oxides were calcinated in quartz crucibles at 1300 K for 5 h in the air to remove surface adsorbates and decompose the hydrocarbonates commonly present on the rare-earth oxide surface. The acid solutions were measured by means of glass measuring graduates with the accuracy of 0.1 ml.

For the synthesis of BaPrCuS<sub>3</sub>, the starting reagents were taken at the stoichiometric molar ratio 6Cu:  $Pr_6O_{11}$ : 6BaCO<sub>3</sub> (Cu – 0.8151 g,  $Pr_6O_{11}$  – 2.1866 g, BaCO<sub>3</sub> – 2.5289 g). For BaSmCuS<sub>3</sub>, the starting reagents were taken at the stoichiometric molar ratio 2Cu:Sm<sub>2</sub>O<sub>3</sub>:2BaCO<sub>3</sub> (Cu – 0.7967 g, Sm<sub>2</sub>O<sub>3</sub> – 2.1832 g, BaCO<sub>3</sub> – 2.4718 g). The charge was inserted in a heat-resistant flat-bottom beaker and, then, concentrated HNO<sub>3</sub> 10 ml (~21% excess) was added. For the best dissolution, the mixture was warmed up on a laboratory tile for the total disappearance of solid components. The solutions of initial nitrates were formed by reactions:

$$\begin{split} 6BaCO_3 + Pr_6O_{11} + 6Cu + 54HNO_3 &\rightarrow 6Ba(NO_3)_2 + 6Pr(NO_3)_3 + 6CuNO_3 + 27H_2O + 12NO_2\uparrow + 6CO_2\uparrow + 4O_2\uparrow \end{split}$$

 $2BaCO_3 + Sm_2O_3 + 2Cu + 18HNO_3 \rightarrow 2Ba(NO_3)_2 + 2Sm(NO_3)_3 + 2CuNO_3 + 9H_2O + 4NO_2\uparrow + 2CO_2\uparrow$ 

Then, the temperature was increased to 500-600 K and the solution was evaporated up to the formation of a solid. On the heating, the solution color changed from blue to green and the brown gas (NO<sub>2</sub>) appeared. That indicates the formation of a compound or a complex on the basis of Cu. The high temperature was kept up to the evaporation of all volatile components. As a result of this treatment, a black color cake was formed. After the evaporation stage, the beaker was cooled and the cake was extracted. Then, the solid cake was ground in an agate mortar. The powder was inserted in a fuzzed quartz boat and treated in the muffle furnace at 1170 K until the end of residual brown gas extraction. After this moment, the sample was calcinated at 1170 K for 2 h in the air. After the annealing, the quartz boat was cooled to room temperature and the final product was ground in an agate mortar. As a result, the stoichiometric oxide mixture was formed.

The sulphidation process was carried out in a vertical quartz reactor at the temperature of 1270 K. The oxide charge was placed in a silica container and inserted in the reactor. Then the reactor was heated to 1270 K and filled with a mixture of argon and sulfiding gas  $CS_2/H_2S$  generated by the decomposition of ammonium thiocyanate NH<sub>4</sub>CSN [33]. For each substance, the reaction was conducted for 5 h:

$$2BaO + Ln_2O_3 + Cu_2O + 3CS_2 \rightarrow 2BaLnCuS_3 + 3CO_2$$

The final powder products of BaLnCuS<sub>3</sub> were of flavovirent color, as shown in Figure 1. As it appears, the color similarity of the compounds is dominated by the presence of  $Cu^+$ . The following advantages of this synthesis route could be mentioned:

- Sedimentation from the nitrate solution allows for the sample homogenization that reduces the time needed for the synthesis of oxide precursor and increases the final product yield.

- Less time needed for the synthesis of sulfide compound in comparison with that typical of ampule synthesis methods.

- Oxide precursors stable at normal conditions are used in the sulfide synthesis.

**Characterization**. The powder diffraction data of BaLnCuS<sub>3</sub> for Rietveld analysis were collected at room temperature with a Bruker D8 ADVANCE powder diffractometer (Cu-Ka radiation) and linear VANTEC detector. The step size of 20 was 0.016°. The 20 range of 10-70° was measured with a 0.6 mm divergence slit and the counting time was 7 s per step, but the 20 range of 70-150° was measured with a 2 mm divergence slit and 12 sec per step. Larger slits allow for a noticeable increase of the intensity of high-angle peaks without a loss of resolution because the high-angle peaks are broad enough to be not affected by a bigger divergence beam. The esd's  $\sigma(I_i)$  of all points on the patterns were calculated using intensities  $I_i$ :  $\sigma(I_i)=I_i^{1/2}$ . The intensities and obtained esd's were further normalized:  $I_i$  norm =  $I_i \times 0.6 \times 7/(\text{slit width} \times \text{time per step})$ ,  $\sigma_{\text{norm}}(I_i) = \sigma(I_i) \times 0.6 \times 7/(\text{slit width} \times \text{time per step})$ , taking into account the actual divergence slit width value, which was used to measure each particular intensity  $I_i$ , and saved in the xye-type file. So, the transformed powder pattern has a usual view in the whole 20 range of 5-150°, but all high-angle points have small esd values.

The optical absorption measurements were carried out for powder samples using a UV-2600 Series Shimadzu spectrometer. The UV-visible spectra were recorded over the wavelength range of 190-1400 nm at room temperature. The particle micromorphology was observed by SEM using a LEO 1430 device. The Raman spectra were collected in backscattering geometry using a triple Raman spectrometer Jobin Yvon T64000 operating in double subtractive mode. The spectra were detected by a liquid nitrogen-cooled CCD cooled to 140 K. The spectral resolution for the recorded Stokes side Raman spectra was about 1 cm<sup>-1</sup> (this resolution was achieved by using gratings with

1800 grooves mm<sup>-1</sup> and 100  $\mu$ m slits). The deformation of the low-frequency spectral edge by an optical slit, which sometimes smears the right features of the low-frequency spectra, was carefully eliminated by rigorous optical alignment. Raman spectra of BaPrCuS<sub>3</sub> were recorded at the 514.5 nm excitation wavelength from Spectra Physics Stabilite 2017 laser. To avoid possible excitation of the luminescence of Sm<sup>3+</sup> ion at the transition to <sup>4</sup>F<sub>3/2</sub> in case of the excitation source specified above, the 647.1 nm line of Lexel 95K laser at 100 mW power (5 mW on the sample) was used as an excitation light source for Raman spectroscopy of BaSmCuS<sub>3</sub>.

The simultaneous thermal analysis was performed in the He (99,999%, Russia) flow with use of a STA 449 F3 Jupiter instrument equipped with a (W3%Re – W25%Re) thermocouple. The analyzed powder sample weighed (90-100) $\pm$ 0.01 mg. The temperature adjustment accuracy was not above 0.3 K. In the temperature range, where thermal events were observed, the heating rate was 10 K/min. The results of DSK/TG experiments were processed in the Proteus-6 programs package [34]. The possible error in the phase transition enthalpy determination was  $\pm$ 18%.

#### 3. Results and discussion

Rietveld refinement was performed using TOPAS 4.2 [35] which accounts the esd's of each point by a special weight scheme. All peaks measured for BaPrCuS<sub>3</sub> were indexed by an orthorhombic unit cell (space group *Pnma*) with cell parameters close to those of BaNdCuS<sub>3</sub> [18]. However, BaSmCuS<sub>3</sub> peaks were indexed by an orthorhombic cell (space group *Cmcm*) with parameters close to those of KThCuS<sub>3</sub> [36]. Besides, in this sample, a small amount (~5% wt.) of BaCu<sub>2</sub>S<sub>2</sub> (space group *Pnma*) [37] impurity was detected. These crystal structures were taken as the starting models for Rietveld refinements.

In order to obtain the model of BaPrCuS<sub>3</sub>, the Nd<sup>3+</sup> site in BaNdCuS<sub>3</sub> was assumed to be occupied by the  $Pr^{3+}$  ion. To obtain the BaSmCuS<sub>3</sub> model, the K<sup>+</sup> site in KThCuS<sub>3</sub> was replaced by Ba, and Th<sup>4+</sup> was replaced by Sm<sup>3+</sup>. The thermal parameters of Ba<sup>2+</sup>, Pr<sup>3+</sup> and Sm<sup>3+</sup> ions were refined anisotropically, but Cu<sup>+</sup> and S<sup>2-</sup> ions were refined with isotropic thermal parameters. The refinements were stable and gave low *R*-factors (Table 1, Figure 2). The obtained coordinates, main bond lengths and anisotropic thermal parameters are in Table 2, Tables 1S and 2S, respectively. The crystal structures of BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub> are presented in Figure 3. Further details of the crystal structures may be obtained from Fachinformationszentrum Karlsruhe, 76344 Eggenstein-Leopoldshafen, Germany (fax: (+49)7247-808-666; E-mail: crystdata@fiz-karlsruhe.de; http://www.fiz-karlsruhe.de/request\_for\_deposited\_data.html on quoting the deposition numbers CSD-1948871 and 1948872.

Thus, two different structures are observed in BaLnCuS<sub>3</sub> compounds for the sequence Sm-Pm-Nd-Pr. On the one hand, BaSmCuS<sub>3</sub> crystallizes in space group *Cmcm*, and, on the other hand, BaNdCuS<sub>3</sub> and BaPrCuS<sub>3</sub> crystallize in space group *Pnma*. The BaPmCuS<sub>3</sub> structure remains unknown. Accordingly, it is valuable to define the relation governing the structure type in ABCX<sub>3</sub> crystals. In order to make the space group prediction for the compounds with general formula ABCX<sub>3</sub>, it was decided to consider the structures known for X = S, Se from the ICSD database. Generally, the structure consists of AX<sub>7</sub>, BX<sub>7</sub> and CX<sub>4</sub> polyhedra columns parallel to the *b*-axis, and they join each other by edges, as shown in Figure S1. The crystal structure deformation and the change of the corresponding space group should depend on ion radii IR(A), IR(B), IR(C) because the column lengths directly depend on them. As it can be reasonably assumed, anion ion radii IR(X)should lead to the overall cell volume deformation without a drastic change of a structure and a space group should not depend on this factor. It is valuable to find the tolerance factor t(IR(A),IR(B),IR(C)), which controls the space group. But the function is not clear. To solve this problem, a deep learning model with three input neurons IR(A), IR(B), IR(C), two intermediate layers with two neurons and one output neuron was built in package Python 3.6 [38]. The conventional sigmoid activation function was used for all layers. In the model, the output neuron with the 0 value means space group *Pnma*, but 1 means the *Cmcm* group. About 2000 epochs of learning were launched and, finally, the small loss function of 5.2% was reached with a high accuracy. The first weight matrix revealed that weights associated with IR(A), IR(C) neurons are

big and positive in the range of 1.3-2. The weights associated with IR(B) are also big, but negative, and the weight values are in the range from -3.4 to -3.0. Therefore, we can conclude that the tolerance factor should proportionally depend on IR(A) and IR(C) values and in the inverse ratio on IR(B). Moreover, it was found that weights corresponding to the IR(X) neuron are small, and it is in the range of 0.3-0.4. One can really conclude that the output neuron value, i.e. space group, almost does not depend on IR(X). In addition, as seen in Table 3S [39-54], the pair compounds ABCX<sub>3</sub> with X = S and X = Se are always crystallized in the same space group. Finally, the dimensionless tolerance factor  $t = IR(A) \times IR(C)/IR(B)^2$  was suggested. All ABCX<sub>3</sub> (X=S, Se) compounds from ICSD were sorted by increasing the *t* factor and are presented in Table 3S, and the critical value t = 0.908 defines the demarcation line between the *Cmcm* and *Pnma* space groups. As shown in Figure 4, all ABCX<sub>3</sub> (X=S, Se) compounds can be presented as points in the field defined with two parameters IR(A)\*IR(C) and IR(B)<sup>2</sup>, and the line  $t = IR(A) \times IR(C)/IR(B)^2 = 0.908$  cuts the space into two fields with different space groups. Only two compounds PbYCuS<sub>3</sub> and SrYCuSe<sub>3</sub> (or PbTbCuS<sub>3</sub> and PbDyCuS<sub>3</sub>) positioned in the vicinity to the demarcation line and the rekated points are away from the general trend. First, the structures of these compounds should be investigated more carefully. As for the prediction of the structure type for BaPmCuS<sub>3</sub>, it seems to be impossible because the related point is assumed to appear on the demarcation line.

The thermal, morphological and spectroscopic properties were measured only for phase-pure BaPrCuS<sub>3</sub>. As it was shown by the XRD analysis, the available BaSmCuS<sub>3</sub> sample contains a significant impurity of BaCu<sub>2</sub>S<sub>2</sub>. Several attempts were made in the synthesis of phase-pure BaSmCuS<sub>3</sub> by the sulphidation technique with some variation of the technological parameters. However, in all experiments, the presence of BaCu<sub>2</sub>S<sub>2</sub> was observed in the final products. The source of this effect is presently unknown, and, accordingly, the synthesis conditions should be optimized to get the phase-pure BaSmCuS<sub>3</sub> product.

The BaPrCuS<sub>3</sub> powder micromorphology is shown in Figure 5. The powder contains irregular partly agglomerated particles of  $\sim$ 1-5 µm in size. The faceted shapes were not detected and

this may be a result of a drastic anion exchange during the sulpfidation reaction at high temperatures. Under the conditions, the formation of faceted microcrystal shapes seems to be highly improbable. The thermal properties of BaPrCuS<sub>3</sub> powder were analyzed over the temperature range from 300 to 1800 K and the related curves are shown in Figure 6. The sample weight was 102.8 mg and the weight loss was not registered in the temperature cycling. The endothermic effect related to the substance melting was observed at  $T_{melt} = 1580.9$  K on heating and the exothermic effect due to the substance crystallization was detected at  $T_{cryst} = 1607.2$  K. Thus, the effect of overcooling is expressed for BaPrCuS<sub>3</sub> and the difference between melting and crystallization points is 27 K. The effect of overcooling is bound to the incongruent type of melting. On the heating curve, the melting peak splitting into two components at 1599 K and 1612 K is evident. The heat of melting determined for BaPrCuS<sub>3</sub> is equal to  $\Delta H_{melt} = 86.6$  kJ·mol<sup>-1</sup>.

At present, the thermal characteristics of ALnCuS<sub>3</sub> compounds are less studied. The thermal parameters available in the literature are summarized in Table 3 [55]. It should be noted that nearly all observed ALnCuS<sub>3</sub> compounds melt incongruently. As seen, the melting temperature values obtained for SrLnCuS<sub>3</sub> and EuLnCuS<sub>3</sub> sulfides are similar for the same Ln element. It is particularly interesting to compare the melting temperatures of the APrCuS<sub>3</sub> (A = Sr, Eu, Ba) set. The melting temperature of BaPrCuS<sub>3</sub> is noticeably higher than that of APrCuS<sub>3</sub> (A = Sr, Eu) and, presumably, this relation indicates a higher thermal stability of the sulfides with bigger radius of the A element. The effect of lanthanide compression can also increase the thermal stability of ALnCuS<sub>3</sub> compounds. Indeed, the BaGdCuS<sub>3</sub> melts congruently at as high temperature as 1685 K [22] and the temperature is significantly higher than that of APrCuS<sub>3</sub> (A = Sr, Eu) and, it is similar to the values known for simple sulfides Ln<sub>2</sub>S<sub>3</sub> [56].

To the best of our knowledge, the calculations of the band structures of BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub> are absent in the literature. However, in this situation, the available band structure calculations for monoclinic sulfides Ba<sub>2</sub>Cu<sub>2</sub>ThS<sub>5</sub> and Ba<sub>2</sub>Cu<sub>2</sub>US<sub>5</sub> (space group C2/m) [57] can be

used for the initial observation of the BaPrCuS<sub>3</sub> band structure. Thus, it is reasonably assumed that the valence band top in BaPrCuS<sub>3</sub> could formed by S 2p, Cu 3d and Pr 4f orbitals, while the conduction band bottom is formed by the contributions from all ions, including 5d orbitals of Ba and Pr. The strong dependence of the band energy on the direction in Brillouin zone is expected and, therefore, both direct and indirect allowed transitions can contribute to the formation of the fundamental absorption edge.

The Tauc plot, in the case of indirect allowed transitions, for the UV-visible spectra recorded for BaPrCuS<sub>3</sub>, is shown in Figure 7. The optical bandgap estimated for BaPrCuS<sub>3</sub> is equal to 2.1 eV. This value is close to the bandgap value 2.00 eV earlier determined for BaLaCuS<sub>3</sub> [17], as well as to the bandgap of another recently studied sulfide EuErCuS<sub>3</sub> (1.94 eV) [58]. At the same time, the imperfect linearity of experimental spectrum in the region up to 2.5 eV indicates the contribution of other types of transitions in accordance with the previously formulated assumption. Besides the fundamental absorption band, the minor narrow peak at ~2 eV can be detected in the spectrum and it is ascribed to the *f*-*f* transition of Pr<sup>3+</sup> ions, namely, <sup>3</sup>H<sub>4</sub> -<sup>1</sup>D<sub>2</sub>. Other prominent *f*-*f* transitions of the Pr<sup>3+</sup> ion, namely, <sup>3</sup>H<sub>4</sub>-<sup>3</sup>P<sub>1</sub> (J=0,1,2) and <sup>1</sup>S<sub>0</sub>, cannot be resolved in the obtained spectrum, and, therefore, it should be deduced that the <sup>3</sup>P<sub>1</sub> levels of Pr<sup>3+</sup> ions should be supposed as contributing to the formation of fundamental absorption band.

The emission spectrum obtained from BaPrCuS<sub>3</sub> under the excitation at 514.5 nm is shown in Figure S2. The left part of this spectrum in the range below 525 nm is formed by the Raman spectrum with its narrow lines. At the wavelengths above 525 nm, the broad band peaking at 650 nm is observed. This spectral band must be assigned to the luminescence from the conduction band of the electronic structure of BaPrCuS<sub>3</sub>. The contribution from the *f*-*f* transitions of the Pr<sup>3+</sup> ion to the luminescence is obviously absent since the 514.5 nm excitation is inefficient to excite the energy level system of *f* electrons in the Pr<sup>3+</sup> ion. The luminescence intensity is comparable to that of Raman spectrum. Therefore, BaPrCuS<sub>3</sub> is highly likely to be an indirect wide-bandgap semiconductor, in accordance with the assumption stated above. The energy spacing of spectral replicas observed in the luminescent band is equal to 0.05 eV (400 cm<sup>-1</sup>). The corresponding vibrational frequencies are absent in the crystal structure of BaPrCuS<sub>3</sub>. Therefore, the origin of these replicas cannot be ascribed to the well-known phonon replicas in the exciton luminescence and needs a special investigation.

The Raman spectra from BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub> powders are given in Figure 8. For BaPrCuS<sub>3</sub>, belonging to space group *Pnma*, the mechanical representation at the Brillouin zone (BZ) center can be written as  $\Gamma_{vibr} = 12A_g + 6A_u + 6B_{1g} + 12B_{1u} + 12B_{2g} + 6B_{2u} + 6B_{3g} + 12B_{3u}$ , where the Raman active modes are  $\Gamma_{Raman} = 12A_g + 6B_{1g} + 12B_{2g} + 6B_{3g}$  [59]. The vibrational irreducible representation at the  $\Gamma$  point of BZ for BaSmCuS<sub>3</sub> is  $\Gamma_{vibr} = 5A_g + 2A_u + 4B_{1g} + 7B_{1u} + B_{2g} + 7B_{2u} + 5B_{3g} + 5B_{3u}$ , where Raman active modes are  $\Gamma_{Raman} = 5A_g + 4B_{1g} + B_{2g} + 5B_{3g}$ . The difference between the Cu–S1 (Cu–S3) bond lengths in BaPrCuS<sub>3</sub> and Cu–S2 in BaSmCuS<sub>3</sub> is equal to 0.007 Å. Thus, the positions of spectral bands, related to the vibrational modes of CuS<sub>4</sub> tetrahedron, should not differ significantly, but the number of spectral bands should not be the same. The set of Raman spectra in several scattering geometries should be obtained to get the bands related to specified Raman modes. However, in the case of powder samples, such measurements are impossible.

In order to interpret the observed Raman spectra, the calculation of BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub> lattice dynamics using the program package LADY [60] was performed. The atomic vibration values were obtained using the simplified version of Born-Karman model [61]. Within this model, only the pair-wise interactions and bond-stretching force constants *F* are considered and the model implies that *F* depends on  $r_{ij}$  (interatomic distance). The  $F(r_{ij})$  dependences are the same for all atom pairs:  $F(r_{ij}) = \lambda \exp(-r_{ij}/\rho)$ , where  $\lambda$  and  $\rho$  are the parameters characterizing a selected pair interaction. To find the model parameters, a special optimization program was written and tested for several representative compounds [62-68]. As it was noted above, BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub> are members of the ALnCuS<sub>3</sub> sulfide family and their structures are formed by the layers of CuS<sub>4</sub> tetrahedra and LnS<sub>6</sub> octahedra [29,30]. According to the lattice dynamics simulations, vibrations at low wavenumbers (20-115 cm<sup>-1</sup>) are related to the vibration of such layers. For example, the lowest wavenumber Raman-active mode in BaPrCuS<sub>3</sub> is graphically shown in Figure 9(a). The strong bands at 62 and 67 cm<sup>-1</sup> in BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub> are related to the antisymmetric displacements of structural layers, and the vibrations are shown in Figures 9(b) and 9(c), respectively. The difference between BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub> in Raman spectra should be in the range of 115-185 cm<sup>-1</sup>. There are no trace of Raman-active vibrational modes in this range in the case of BaSmCuS<sub>3</sub>, but a group of lines is observed in the case of BaPrCuS<sub>3</sub> (Figure 8). The region of Raman spectra between 188 and 250 cm<sup>-1</sup> is related to the vibrations of sulfur ions mainly and, for example, the bending vibration of CuS<sub>4</sub> in BaSmCuS<sub>3</sub> is shown in Figure 9(d). It is clearly seen in Figure S1 that CuS<sub>4</sub> tetrahedra from neighboring unit cells share one common vertex and the related vibration of bringing sulfur ion appears in the spectra as a spectral band around 242 cm<sup>-1</sup> (Figure 9(e)). The stretching-like vibrations in CuS<sub>4</sub> tetrahedra are in the high wavenumber range of 310-350 cm<sup>-1</sup>, and they are shown in Figure 9(f).

#### 4. Conclusion

This study addresses the synthesis, structure, optical and thermal properties of the new complex sulfides BaPrCuS<sub>3</sub> (space group *Pnma*) and BaSmCuS<sub>3</sub> (space group *Cmcm*). The compounds crystallize in two different orthorhombic structures that open a possibility to analyze the mechanism of the structure transformation in the known ABCX<sub>3</sub> (X = S, Se) compounds for the sequence Sm-Pm-Nd-Pr. To solve this problem, the deep learning neuron model was built on the base of the known ABCX<sub>3</sub> compounds. The introduced tolerance factor  $t = IR(A) \times IR(C)/IR(B)^2$  can be used as a robust indicator for the prediction of the structure type for compounds ABCX<sub>3</sub> (X = S, Se), which crystal structure remains unknown. BaPrCuS<sub>3</sub> is shown to be the indirect bandgap compound with bandgap width 2.1 eV.

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Compound	BaPrCuS <sub>3</sub>	BaSmCuS <sub>3</sub>
Sp. Gr.	Pnma	Стст
<i>a</i> , Å	10.56074 (6)	4.07269 (4)
b, Å	4.11305 (2)	13.4499 (1)
c, Å	13.42845 (7)	10.3704 (1)
V, Å <sup>3</sup>	583.289 (5)	568.06 (1)
Ζ	2	2
2θ-interval, °	10-150	10-150
$R_{wp},$ %	1.18	1.84
$R_{p}, \%$	1.20	1.39
$R_{exp}$ , %	0.62	0.86
$\chi^2$	1.89	2.14
$R_B, \%$	0.79	0.88

Table 1. Main parameters of processing and refinement of the BaPrCuS3 and BaSmCuS3 samples

**Table 2.** Fractional atomic coordinates and isotropic or equivalent isotropic displacement parameters  $(Å^2)$  of BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub>

	x	у	Z	$U_{ m iso}$ */ $U_{ m eq}$
		BaPrCuS <sub>3</sub>		
Pr	0.00804 (8)	0.25	0.75733 (6)	0.0041 (4)
Cu	0.2406 (2)	0.25	0.28297 (9)	0.0093 (4)*
Ba	0.26735 (8)	0.25	0.00560 (4)	0.0178 (5)
<b>S</b> 1	0.0555 (2)	0.25	0.3822 (2)	0.0038 (6)*
<b>S</b> 2	0.2558 (3)	0.25	0.6825 (1)	0.0039 (5)*
<b>S</b> 3	0.4181 (2)	0.25	0.3905 (2)	0.0070 (7)*
		BaSmCuS	3	
Sm	0.5	0	0	0.71 (3)
Cu	0.5	0.53447 (15)	0.25	1.01 (6)*
Ba	0.5	0.25634 (7)	0.25	1.07 (12)
<b>S</b> 1	0.5	0.63449 (19)	0.0670 (2)	0.85 (7)*
S2	0.5	0.9328 (2)	0.25	1.10 (10)*

Compound	T <sub>melt</sub> , K	$H_{melt}$ , kJ·mol <sup>-1</sup>	Reference
SrLaCuS <sub>3</sub>	1513	6,9	24,25
SrCeCuS <sub>3</sub>	1468	5,2	25,46
SrPrCuS <sub>3</sub>	1459	13,2	25
SrNdCuS <sub>3</sub>	1429	16,8	25
SrSmCuS <sub>3</sub>	1605	2,8	25
EuLaCuS <sub>3</sub>	1539	23,4	32, 55
EuCeCuS <sub>3</sub>	1524	23,7	55
EuPrCuS <sub>3</sub>	1497	20,1	55
EuNdCuS <sub>3</sub>	1470	17,7	39,55
EuDyCuS <sub>3</sub>	1727	2,9	28
BaGdCuS <sub>3</sub>	1685	-	39
BaPrCuS <sub>3</sub>	1580.9	86.6	This study

**Table 3.** Temperature of melting and enthalpy of melting in compounds  $ABCuS_3$ 

# Captions

Figure 1. Photo of (a) BaPrCuS<sub>3</sub> and (b) BaSmCuS<sub>3</sub> samples.

Figure 2. Difference Rietveld plots of (a) BaPrCuS<sub>3</sub> and (b) BaSmCuS<sub>3</sub>.

Figure 3. Crystal structures of (a)  $BaPrCuS_3$  and (b)  $BaSmCuS_3$  sulfides. The anisotropic thermal parameters of  $Pr^{3+}$ ,  $Sm^{3+}$  and  $Ba^{2+}$  ions are shown at the 50% probability level.

Figure 4. Structure types in known ABCX<sub>3</sub> (X = S, Se) crystals.

Figure 5. SEM image of BaPrCuS<sub>3</sub> powder.

Figure 6. DTA curves recorded for BaPrCuS<sub>3</sub>. The baselines are shown in red color.

Figure 7. Tauc plot for BaPrCuS<sub>3</sub>.

Figure 8. Raman spectra of BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub>. The black, red, blue and green vertical ticks mark the calculated wavenumbers of Raman-active vibrational modes.

Figure 9. Examples of calculated vibrational modes in BaPrCuS<sub>3</sub>, (a)  $A_g$  20.6 cm<sup>-1</sup> low frequency layer vibration, (e)  $A_g$  202.5 cm<sup>-1</sup> related to the vibration of bridging oxygen atoms and, in BaSmCuS<sub>3</sub>, (b) and (c)  $A_g$  64.6 cm<sup>-1</sup> and  $B_{3g}$  68.6 cm<sup>-1</sup> layer vibrations, respectively, (d)  $A_g$  209.2 cm<sup>-1</sup> O-S–O bending-like vibration, (f)  $A_g$  330.6 cm<sup>-1</sup> S-O stretching-like vibration.





Figure 2



b

Figure 3



Figure 4



Figure 5



Figure 6.



Figure 7



Figure 8



Fig. 9

#### **Supporting information**

Synthesis, structural and spectroscopic properties of orthorhombic compounds BaLnCuS3 (Ln = Pr, Sm)

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BaPrCuS <sub>3</sub>				
Pr—S1 <sup>i</sup>	2.862 (2)	Cu—S3	2.366 (3)	
Pr—S2	2.803 (3)	Ba—S1 <sup>iv</sup>	3.237 (2)	
Pr—S2 <sup>ii</sup>	2.784 (3)	Ba—S1 <sup>v</sup>	3.396 (3)	
Pr—S3 <sup>iii</sup>	2.835 (2)	Ba—S2 <sup>iv</sup>	3.151 (2)	
Cu—S1	2.366 (3)	Ba—S3 <sup>iv</sup>	3.233 (2)	
Cu—S2 <sup>iv</sup>	2.460 (1)			
BaSmCuS <sub>3</sub>				
Sm—S1 <sup>i</sup>	2.8110 (17)	Cu—S2 <sup>iii</sup>	2.453 (2)	
Sm—S2 <sup>ii</sup>	2.7457 (11)	Ba—S1 <sup>iii</sup>	3.230 (2)	
Cu—S1	2.326 (3)	Ba—S2 <sup>iii</sup>	3.127 (3)	

Table 1S. Main bond lengths (Å) of BaPrCuS<sub>3</sub> and BaSmCuS<sub>3</sub>

Symmetry codes:

BaPrCuS<sub>3</sub>: (i) -*x*, -*y*, -*z*+1; (ii) *x*-1/2, -*y*+1/2, -*z*+3/2; (iii) -*x*+1/2, -*y*, *z*+1/2; (iv) -*x*+1/2, -*y*, *z*-1/2; (v) *x*+1/2, -*y*+1/2, -*y*+1/2, -*z*+1/2; (iv) -*x*+1/2, -*y*, *z*-1/2; (v) *x*+1/2, -*y*+1/2, -*y*+1/2, -*y*+1/2, -*y*+1/2; (iv) -*x*+1/2, -*y*, *z*-1/2; (v) *x*+1/2, -*y*+1/2, -*y*+1/2, -*y*+1/2; (iv) -*x*+1/2, -*y*, *z*-1/2; (v) *x*+1/2, -*y*+1/2, -*y*+1/2, -*y*+1/2; (iv) -*x*+1/2, -*y*+1/2; (iv) -*x*+1/2, -*y*+1/2; (iv) -*x*+1/2, -*y*+1/2; (iv) -*x*+1/2; (v) -*x*+1/2;

BaSmCuS<sub>3</sub>: (i) -*x*+1/2, -*y*+1/2, -*z*; (ii) *x*, *y*-1, *z*; (iii) -*x*+1/2, *y*-1/2, -*z*+1/2

	$U^{11}$	$U^{22}$	$U^{33}$	$U^{12}$	$U^{13}$	$U^{23}$
BaPrCuS <sub>3</sub>						
Pr	0.0041 (4)	0.0039 (3)	0.0058 (4)	0	-0.0003 (3)	0
Ba	0.0178 (5)	0.0021 (4)	0.0036 (4)	0	-0.0010 (3)	0
BaSmCuS <sub>3</sub>						
Sm	0.0109 (7)	0.0100 (8)	0.0062 (8)	0	-0.0060 (6)	0
Ba	0.0061 (8)	0.0086 (9)	0.0257 (1)	0	0	0

**Table 2S.** Atomic displacement parameters (Å<sup>2</sup>)

**Table 3S.** Chemical formulas, space groups and corresponding tolerance factors of ABCX3

 compounds

Compound	$t = IR(A)*IR(C)/IR(B)^2$	Space group	Reference
EuNdCuS <sub>3</sub>	0.704209	Pnma	[39]
SrNdCuS <sub>3</sub>	0.709843	Pnma	[40]
PbBiCuS <sub>3</sub>	0.729569	Pnma	[41]
SrLaCuSe <sub>3</sub>	0.741104	Pnma	[42]
SrLaCuS <sub>3</sub>	0.741104	Pnma	[24]
EuPrCuS <sub>3</sub>	0.765228	Pnma	[43]
SrPrCuS <sub>3</sub>	0.77135	Pnma	[44]
SrPrCuSe <sub>3</sub>	0.77135	Pnma	[45]
SrCeCuS <sub>3</sub>	0.782375	Pnma	[46]
SrCeCuSe <sub>3</sub>	0.782375	Pnma	[45]
BaNdCuS <sub>3</sub>	0.799982	Pnma	[18]
EuSmCuS <sub>3</sub>	0.817204	Pnma	[47]
BaLaCuS <sub>3</sub>	0.835212	Pnma	[17]
BaLaCuSe <sub>3</sub>	0.835212	Pnma	[17]
Eu <sub>2</sub> CuS <sub>3</sub>	0.836298	Pnma	[48]
SrGdCuS <sub>3</sub>	0.859243	Pnma	[44]
SrGdCuSe <sub>3</sub>	0.859243	Pnma	[42]
PbGdCuSe <sub>3</sub>	0.879701	Pnma	[49]
BaCeCuS <sub>3</sub>	0.881724	Pnma	[18]
BaCeCuSe <sub>3</sub>	0.881724	Pnma	[18]
EuDyCuS <sub>3</sub>	0.90172	Pnma	[50]
Y <sub>0.11</sub> Yb <sub>0.89</sub> YbCuS <sub>3</sub>	0.907855	Pnma	[51]
PbTbCuS <sub>3</sub>	0.908526	Стст	[49]
PbDyCuS <sub>3</sub>	0.930575	Стст	[49]
SrYCuSe <sub>3</sub>	0.933333	Pnma	[52]
PbHoCuS <sub>3</sub>	0.953436	Стст	[49]
PbYCuS <sub>3</sub>	0.955556	Pnma	[53]
BaGdCuS <sub>3</sub>	0.968353	Стст	[18]
BaGdCuSe <sub>3</sub>	0.968353	Стст	[54]
BaLaCuS <sub>3</sub>	0.977149	Стст	[17]
PbTmCuS <sub>3</sub>	0.999483	Стст	[14]

SrLuCuS <sub>3</sub>	1.019801	Стст	[44]
SrLuCuSe <sub>3</sub>	1.019801	Стст	[42]
PbYbCuS <sub>3</sub>	1.02731	Стст	[14]
PbLuCuS <sub>3</sub>	1.044082	Стст	[14]
BaErCuS <sub>3</sub>	1.075622	Стст	[18]
BaErCuSe <sub>3</sub>	1.075622	Стст	[18]
BaLaAgSe <sub>3</sub>	1.39202	Стст	[18]
BaErAgSe <sub>3</sub>	1.792703	Стст	[18]
BaGdAuSe <sub>3</sub>	2.001264	Стст	[19]



**Figure S1.** Three columns of polyhedra: AX<sub>7</sub>, BX<sub>7</sub> and CX<sub>4</sub> which are linked with each other by edges.



Figure S2. Emission spectrum of BaPrCuS<sub>3</sub> recorded under the excitation at 514.5 nm.