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# On One Two-dimensional Stationary Flow of a Binary Mixture and Viscous Fluid in a Plane Layer 

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#### Abstract

Received 12.10.2015, received in revised form 10.11.2015, accepted 20.12.2015 Nonlinear model of convection in Oberbeck-Boussinesq approximation describing the flat joint motion of a binary mixture and viscous fluid with a common interface is investigated. It is important that the longitudinal temperature gradient and the concentration is quadratic dependence on the coordinate $x$. Stationary solution of the system is built.


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Convection is one of the most common of hydrodynamic phenomena in nature. Study of convection is an important part of the theoretical fluid mechanics. Natural convective motions occur in an inhomogeneous field of mass forces caused by the ununiform heating of the liquid. The area of practical applications of this phenomenon is very wide. Convective processes influence the thermal conditions in the oil storage tanks, chemical process technology and others. Theoretical study of natural convection usually deals with equations of motion in the OberbeckBoussinesq approximation. Problems for thermal convection are very complex because of the diversity of cavities and thermal boundary conditions for the nonlinear system of partial differential equations. Solutions of the Oberbeck-Boussinesq equations with a linear dependence of temperature on one of the space coordinates firstly were studied by G. A. Ostroumov [1]. The exact solution described plane stationary flow in a strip under action of longitudinal temperature gradient and transversal gravity field, was obtained by R. V. Birikh [2]. Some generalizations of this solution taking into account concentration of liquid mixture are described in [3]. The existence of solutions with nonlinear dependence of density on temperature and concentration is proved in [4] where two boundary value problems with exponential temperature distribution on the walls are solved. In [5-7] the exact solutions of the three-dimensional convection problem for two immiscible viscous, incompressible fluids in a channel with a rectangular cross section in the presence of the interface and under the influence of a longitudinal temperature gradient are studied. In this paper we are built the exact solution of the two-layer convection system with longitudinal temperature gradient at the solid walls and shear force of gravity.

## 1. Problem formulation

Let us consider the joint motion of a binary mixture and viscous fluid with a general interface. Suppose that $\Omega_{1}=\left\{|x|<\infty, 0<y<l_{1}\right\}$ is the region occupied by a binary mixture and

[^0]$\Omega_{2}=\left\{|x|<\infty, l_{1}<y<l_{2}\right\}$ is the region with a viscous fluid. The system is bounded by solid walls $y=0$ and $y=l_{2}$ with given temperature distribution on them.

For description of the motion of both region $\Omega_{j}(j=1,2)$ we use the Boussinesq approximation. We assume that the temperature and the concentration differ only slightly from constant mean values therefore the Oberbeck-Boussinesq approximation is valid

$$
\rho_{j}=\rho_{0 j}\left(1-\beta_{j}^{\theta} \theta-\beta_{j}^{c} c\right)
$$

where $\rho_{0 j}$ is the characteristic medium density corresponding to the mean values of the temperature and concentration in the layer $j, \theta$ and $c$ are the deviations from their mean values $(c$ corresponds to the light component), $\beta_{j}^{\theta}$ and $\beta_{j}^{c}$ are the temperature and concentration expansion coefficient; $\beta_{2}^{c}=0$. Then the equations of binary mixture convection can be written in the form

$$
\begin{gather*}
\mathbf{u}_{j t}+\left(\mathbf{u}_{j} \cdot \nabla\right) \mathbf{u}_{j}=-\frac{1}{\rho_{0 j}} \nabla p_{j}+\nu_{j} \Delta \mathbf{u}_{j}-\mathbf{g}\left(\beta_{j}^{\theta}\left(\theta_{j}-\theta_{0 j}\right)+\beta_{j}^{c}\left(c_{j}-c_{0 j}\right)\right) \\
\theta_{j t}+\mathbf{u}_{j} \cdot \nabla \theta_{j}=\chi_{j} \Delta \theta_{j}  \tag{1}\\
c_{1 t}+\mathbf{u}_{1} \cdot \nabla c_{1}=d_{1} \Delta c_{1}+\alpha d_{1} \Delta \theta_{1} \\
\operatorname{div} \mathbf{u}_{j}=0
\end{gather*}
$$

Here $\mathbf{u}_{j}$ is the velocity field, $p_{j}$ is the pressure measured from the hydrostatic pressure corresponding to $p_{0 j}, \rho_{0 j}$ is the density, $\nu_{j}$ is kinematic viscosity, $\chi_{j}$ is the temperature conductivity, $d_{1}$ is the diffusivity, $\alpha d_{1}$ is the thermal diffusion coefficient. Normal thermal diffusion corresponds to the value of $\alpha<0$, and for the anomalous $\alpha>0$.

We introduce the coordinate system with the $x$ axis aligned with the lower boundary of the layer 1 and the $y$ axis directed vertically upward (Fig. 1).


Fig. 1. The scheme of two-layer flow between the rigid walls with interface $y=l_{1}$

Let us define the boundary conditions. On solid walls are put no-slip conditions and the temperature distribution and the absence of mass flux through the walls are written as

$$
\begin{array}{ll}
y=0: & \mathbf{u}_{1}=0, \quad \theta=\theta_{10}(x), \quad c_{1 y}+\alpha \theta_{1 y}=0 \\
y=l_{2}: & \mathbf{u}_{2}=0, \quad \theta=\theta_{20}(x) \tag{2}
\end{array}
$$

At the interface $y=l_{1}$ the conditions of equality of the velocities, kinematic and dynamic conditions are written as

$$
\begin{align*}
& u_{1}=u_{2}, \quad v_{1}=v_{2}=0  \tag{3}\\
& \rho_{2} \nu_{2} u_{2 y}-\rho_{1} \nu_{1} u_{1 y}=-æ_{1} \theta_{1 x}-æ_{2} c_{1 x}
\end{align*}
$$

The condition of temperature continuity and the equality of heat fluxes are as follows

$$
\begin{equation*}
\theta_{1}=\theta_{2}, \quad k_{1} \theta_{1 y}=k_{2} \theta_{2 y} \tag{4}
\end{equation*}
$$

In addition, the condition of absence of mass flux through the interface is

$$
\begin{equation*}
c_{1 y}+\alpha \theta_{1 y}=0 \tag{5}
\end{equation*}
$$

Here $k_{j}=\chi_{j} \rho_{j} c_{p_{j}}$ are the thermal conductivities, $\sigma=\sigma(\theta, c)$ is the coefficient of surface tension. For many mixtures, the linear law provides a good approximation of this dependence

$$
\sigma(\theta, c)=\sigma^{0}-æ_{1}\left(\theta-\theta_{0}\right)-æ_{2}\left(c_{1}-c_{0}\right),
$$

where $æ_{1}>0$ is the temperature coefficient and $æ_{2}$ is the concentration coefficient (usually $æ_{2}<0$ since the surface tension increases with concentration). Constants $\theta_{0}, c_{0}$ are the temperature and concentration values of arbitrary point on interface.

We should to add the initial conditions: $u_{j}=0, \theta_{j}=\theta_{j}^{0}(x, y), c_{1}=c^{0}(x, y)$. All the physical characteristics of the system are assumed to be constant and correspond to the mean temperature and concentration.

## 2. Exact solution of the two-dimensional problem

We find the form of the solution describing the convective flow in the system of liquids with the interface in the form

$$
\begin{gather*}
u_{j}=U_{j}(y, t) x+W_{j}(y, t), \quad v_{j}=V_{j}(y, t) \\
\theta_{j}=A_{j}(y, t) x^{2}+B_{j}(y, t), \quad c_{1}=H_{1}(y, t) x^{2}+E_{1}(y, t)  \tag{6}\\
p_{j}=P(x, y, t)
\end{gather*}
$$

The substitution of solution (6) into equations of motion (1) gives the relations

$$
\begin{gather*}
\nu_{j} U_{j y y}-U_{j t}-U_{j}^{2}-V_{j} U_{j y}=2 g \int_{\Omega_{j}}\left(\beta_{j}^{\theta} A_{j}+\beta_{j}^{c} H_{j}\right) d y+s_{j}(t) \\
\frac{1}{\rho_{j}} P_{j}=\left(\nu_{j} U_{j y y}-U_{j t}-U_{j}^{2}-V_{j} U_{j y}\right) \frac{x^{2}}{2}+h_{j}(y, t)  \tag{7}\\
h_{j y}=\nu_{j} V_{j y y}-V_{j t}-V_{j} V_{j y}+g\left(\beta_{j}^{\theta} B_{j}+\beta_{j}^{c} E_{j}\right) \\
W_{j}=0, \quad V_{j y}=-U_{j}
\end{gather*}
$$

The equations for determining of the temperature and the concentration field take the form

$$
\begin{gather*}
A_{j t}+2 U_{j} A_{j}+V_{j} A_{j y}=\chi_{j} A_{j y y} \\
B_{j t}+V_{j} B_{j y}=\chi_{j}\left(2 A_{j}+B_{j y y}\right)  \tag{8}\\
H_{1 t}+2 U_{1} H_{1}+V_{1} H_{1 y}=d_{1}\left(H_{1 y y}+\alpha_{1} A_{1 y y}\right) \\
E_{1 t}+V_{1} E_{1 y}=d_{1}\left(2 H_{1}+E_{1 y y}+\alpha_{1}\left(2 A_{1}+B_{1 y y}\right)\right) .
\end{gather*}
$$

The following boundary condition at solid walls are held

$$
\begin{gather*}
U_{1}(0, t)=0, \quad U_{2}\left(l_{2}, t\right)=0, \quad A_{1}(0, t)=A_{10}(t), \quad A_{2}\left(l_{2}, t\right)=A_{20}(t) \\
B_{1}(0, t)=B_{10}(t), \quad B_{2}\left(l_{2}, t\right)=B_{20}  \tag{9}\\
H_{1 y}(0, t)+\alpha_{1} A_{1 y}(0, t)=0 ; \quad E_{1 y}(0, t)+\alpha_{1} B_{1 y}(0, t)=0
\end{gather*}
$$

The boundary conditions at the interface $y=l_{1}$ are:

$$
\begin{gather*}
U_{1}=U_{2}, \quad \rho_{2} \nu_{2} U_{2 y}-\rho_{1} \nu_{1} U_{1 y}=-2 æ_{1} A_{1}-2 æ_{2} H_{1}, \\
A_{1}=A_{2}, \quad k_{1} A_{1 y}=k_{2} A_{2 y} ; \\
B_{1}=B_{2}, \quad k_{1} B_{1 y}=k_{2} B_{2 y} ;  \tag{10}\\
H_{1 y}+\alpha_{1} A_{1 y}=0 ; \quad E_{1 y}+\alpha_{1} B_{1 y}=0 ; \\
\int_{0}^{l_{1}} U_{1}(y, t) d y=0, \quad \int_{l_{1}}^{l_{2}} U_{2}(y, t) d y=0 . \tag{11}
\end{gather*}
$$

The first from conditions (11) is a consequence of the kinematic conditions, when the interface is stationary, and the second one is the slip condition for velocity components $V_{2}(y, t)$ on the wall of $y=l_{2}$.

The initial data are written in the form

$$
\begin{gathered}
U_{j}(y, 0)=0, \quad V_{j}(y, 0)=0, \quad A_{j}(y, 0)=a_{j}^{0}(y) \\
B_{j}(y, 0)=b_{j}^{0}(y), \quad H_{1}(y, 0)=H^{0}(y), \quad E_{1}(y, 0)=E^{0}(y)
\end{gathered}
$$

Note that this problem is nonlinear and inverse. Because of the function $s_{j}(t)$ remains unknown as well as functions $U_{j}, V_{j}, A_{j}, B_{j}, H_{1}, E_{1}$.

We introduce the characteristic length scales, time functions $U_{j}, V_{j}, P_{j}, A_{j}, B_{j}, H_{1}, E_{1}, h_{j}, s_{j}$ respectively

$$
\begin{gather*}
x=l_{1} \xi, \quad t=\frac{l_{1}^{2}}{\nu_{1}} \tau, \quad U_{j}=\frac{æ_{1} \triangle A l_{1}}{\rho_{1} \nu_{1}} U_{j}^{*}, \quad V_{j}=\frac{æ_{1} \triangle A l_{1}^{2}}{\rho_{1} \nu_{1}} V_{j}^{*}, \quad P_{j}=æ_{1} \triangle A l_{1} P_{j}^{*}, \\
A_{j}=\triangle A A_{j}^{*}, \quad B_{j}=\triangle A l_{1}^{2} B_{j}^{*}, \quad H_{1}=\frac{\beta_{1}^{\theta} \triangle A}{\beta_{1}^{c}} H_{1}^{*}, \quad E_{1}=\frac{\beta_{1}^{\theta} \triangle A l_{1}^{2}}{\beta_{1}^{c}} E_{1}^{*}  \tag{12}\\
h_{j}=\frac{æ_{1} \triangle A l_{1}}{\rho_{1}} h_{j}^{*}, \quad s_{j}=\frac{æ_{1} \triangle A}{\rho_{1} l_{1}} s_{j}^{*},
\end{gather*}
$$

where $\triangle A=\max _{t \geqslant 0}\left|A_{20}(t)-A_{10}(t)\right|>0$. If $A_{20}(t)=A_{10}(t)$ than $\triangle A=\max _{j} \max _{y}\left|A_{j 0}(y)\right|>0$.
We have the multiplier

$$
\begin{equation*}
M=\frac{æ_{1} \triangle A l_{1}^{3}}{\rho_{1} \nu_{1}^{2}} \tag{13}
\end{equation*}
$$

called the Marangoni number at the nonlinear summands in equations (7)-(8) written with dimensionless variables. Let us mention here also Prandtl number $\operatorname{Pr}_{j}$, Schmidt number $S c_{j}$, parameter $G_{j}=G r_{j} / M$ (where $G r_{j}$ is Grashof number), parameter $\omega$ and split ratio $\psi$ :

$$
P r_{j}=\frac{\nu_{j}}{\chi_{j}}, \quad S c_{j}=\frac{\nu_{j}}{d_{j}}, \quad G_{j}=\frac{g \beta_{j}^{\theta} \rho_{1} l_{1}^{2}}{æ_{1}}, \quad \omega=\frac{æ_{2} \beta_{1}^{\theta}}{æ_{1} \beta_{1}^{c}}, \quad \psi=-\frac{\alpha \beta_{1}^{c}}{\beta_{1}^{\theta}} .
$$

## 3. Stationary solution

In the present section we describe stationary solution of (7)-(8). Assume that the motion is creeping in one layers so $M \ll 1$ and parameter $G_{j}=G r_{j} / M=O(1)$. These conditions can be valid in either thin layer or very viscous fluid according to the formula (13).

In such case the steady state problem (7)-(8) has a special form

$$
\begin{gather*}
\frac{\nu_{j}}{\nu_{1}} U_{j \eta \eta}=2 G_{j} \int_{\Omega_{j}}\left(A_{j}+\frac{\beta_{j}^{c}}{\beta_{1}^{c}} H_{j}\right) d \eta+s_{j} \\
A_{j \eta \eta}=0, \quad B_{j \eta \eta}=-2 A_{j}  \tag{14}\\
H_{1 \eta \eta}=0, \quad E_{1 \eta \eta}=-2 H_{1}
\end{gather*}
$$

$$
\begin{gather*}
V_{j \eta}=-U_{j} \\
\frac{\rho_{1}}{\rho_{j}} P_{j}=\frac{\nu_{j}}{2 \nu_{1}} U_{j \eta \eta} x^{2}+h_{j}  \tag{15}\\
h_{j \eta}=\frac{\nu_{j}}{\nu_{1}} V_{j \eta \eta}+G_{j} B_{j}+\frac{\beta_{j}^{c}}{\beta_{1}^{c}} G_{1} E_{j}
\end{gather*}
$$

Integrating (14) provides a solution to the problem as the

$$
\begin{gather*}
A_{1}=m_{1} \eta+m_{2}, \quad A_{2}=m_{3} \eta+m_{4} \\
B_{1}=-\frac{m_{1}}{3} \eta^{3}-m_{2} \eta^{2}+m_{5} \eta+m_{6}, \quad B_{2}=-\frac{m_{3}}{3} \eta^{3}-m_{4} \eta^{2}+m_{7} \eta+m_{8} \\
H_{1}=m_{9} \eta+m_{10}, \quad E_{1}=-\frac{m_{9}}{3} \eta^{3}-m_{10} \eta^{2}+m_{11} \eta+m_{12}  \tag{16}\\
U_{1}=G_{1}\left(\frac{m_{1}+m_{9}}{12} \eta^{4}+\frac{m_{2}+m_{10}}{3} \eta^{3}\right)+\frac{s_{1}}{2} \eta^{2}+m_{13} \eta+m_{14} \\
U_{2}=\nu G_{2}\left(\frac{m_{3}}{12} \eta^{4}+\frac{m_{4}}{3} \eta^{3}\right)+\frac{\nu s_{2}}{2} \eta^{2}+m_{15} \eta+m_{16}
\end{gather*}
$$

Constants $m_{i}, i=\overline{1,16}, s_{1}, s_{2}$ in (16) determined from the boundary conditions (9)-(11) and have the form

$$
\begin{aligned}
& m_{1}=\frac{l\left(A_{20}-A_{10}\right)}{l-k l+k}, \quad m_{2}=A_{10}, \quad m_{3}=k m_{1}, \quad m_{4}=\frac{A_{20} l-k m_{1}}{l}, \quad m_{6}=B_{10}, \\
& m_{5}=\frac{3 l^{2}(2 k l-2 k-l) A_{10}-3 l(l-1)^{2} A_{20}+3 l^{3}\left(B_{10}-B_{20}\right)+m_{1}\left(l\left(k l^{2}+3 k l-l^{2}-6 k\right)+2 k\right)}{3 l^{2}(k l-k-l)}, \\
& m_{7}=\frac{3 A_{10} k l^{3}+3 l\left(k l^{2}-2 l^{2}-k\right) A_{20}+3 k l^{3}\left(B_{10}-B_{20}\right)+k\left(k l^{3}-3 k l^{2}-l^{3}+6 l^{2}+2 k\right) m_{1}}{3 l^{2}(k l-k-l)}, \\
& m_{8}=\frac{3 A_{10} k l^{2}+3 l(k l-k-2 l+1) A_{20}+3 l^{2}\left(B_{20} l(1-k)-B_{10} k\right)}{3 l^{2}(k l-k-l)}- \\
& -\frac{k\left(k l^{2}-3 k l-l^{2}+2 k+6 l-2\right) m_{1}}{3 l^{2}(k l-k-l)}, \\
& m_{9}=\psi m_{1}, \quad m_{10}=\psi A_{10}, \quad m_{11}=\psi m_{5}, \quad m_{14}=0, \\
& m_{13}=\frac{(\psi+1)\left((7 \rho \nu(l-1)-4 l) m_{1}+5 A_{10}(3 \rho \nu(l-1)-2 l)\right) G_{1}}{60(\rho \nu l-\rho \nu-l)}+ \\
& +\frac{\nu(l-1)^{3}\left(3 k(l-1) m_{1}+5 A_{20} l\right) G_{2}}{60 l^{3}(\rho \nu l-\rho \nu-l)}-\frac{\nu \rho(\psi \omega+1)(l-1)\left(A_{10}+m_{1}\right)}{\rho \nu l-\rho \nu-l}, \\
& m_{15}=\frac{l \rho(\psi+1)(l+2)\left(5 A_{10}+3 m_{1}\right) \nu G_{1}}{60(l \nu \rho-\rho \nu-l)(l-1)}-\frac{l \rho(\psi \omega+1)(l+2)\left(A_{10}+m_{1}\right) \nu}{(l \nu \rho-\rho \nu-l)(l-1)}- \\
& -\frac{5\left(2 \nu(l-1)\left(l^{2}+4 l+1\right) \rho-3 l^{2}(l+3)\right) A_{20} \nu G_{2}}{60 l^{2}(l \nu \rho-\rho \nu-l)}- \\
& -\frac{k m_{1}\left(2 \rho(l-1)\left(2 l^{3}+3 l^{2}-12 l-3\right) \nu-l^{2}\left(7 l^{2}+6 l-33\right)\right) \nu G_{2}}{60 l^{3}(l \nu \rho-\rho \nu-l)},
\end{aligned}
$$

$$
\begin{gathered}
m_{16}=-\frac{\nu \rho(\psi+1)(2 l+1)\left(5 A 10+3 m_{1}\right) G_{1}-60 \nu \rho(\psi \omega+1)(2 l+1)\left(A_{10}+m_{1}\right)}{120(l \nu \rho-\rho \nu-l)(l-1)}+ \\
+\frac{\nu G_{2} A_{20}(l-1)\left(20 l\left(l^{2}-1\right) \rho \nu-5 l\left(6 l^{2}+3 l-1\right)\right)}{120 l^{3}(l \nu \rho-\rho \nu-l)(l-1)}+ \\
+\frac{\nu G_{2}(l-1) k\left(2(l-1)\left(4 l^{2}-3 l-6\right) \rho \nu-14 l^{3}+15 l^{2}+12 l-3\right) m_{1}}{120 l^{3}(l \nu \rho-\rho \nu-l)(l-1)} ; \\
-\frac{\nu(l-1)^{3}\left(3 k(l-1) m_{1}+5 A_{20} l\right) G_{2}-60 l^{3} \nu \rho(\psi \omega+1)(l-1)\left(A_{10}+m_{1}\right)}{20 l^{3}(\rho \nu l-\rho \nu-l)} \\
s_{1}=-\frac{(\psi+1)\left((l-1)\left(25 A_{10}+9 m_{1}\right) \rho \nu-2 l\left(10 A_{10}+3 m_{1}\right)\right) G_{1}}{20(\rho \nu l-\rho \nu-l)}- \\
+\frac{(l-1)^{2}\left(3 k(l-1)(2 \rho \nu(l-1)-3 l) m_{1}+5 A_{20} l(4 \rho \nu(l-1)-5 l)\right) G_{2}}{20(l-1)(l \rho \nu-\rho \nu-l) l^{2}} .
\end{gathered}
$$

The expression for $V_{j}, P_{j}$ can be obtained with help of formulas (15) and have not shown here because cumbersome.

On Fig. 2a we represent the velocity profiles. In the upper layer the vertical velocity component $V$ for parameters $A_{10}=0.1, A_{20}=-0.3, B_{10}=25, B_{20}=20, G_{1}=1.05, G_{2}=0.98$, $M=0.034$ is positive and the horizontal component of the velocity changes sign. The fluid moves vertically when $x=0$ along the $y$ axis and vertically downward in the lower layer and symmetrically rotated about the axis $y$ (Fig. 2b).


Fig. 2. Profiles of the velocity components $U$ and $V(a)$ and isolines of the stream function $\psi_{j}=-V_{j} x(b)$

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# Двумерное стационарное течение бинарной смеси и вязкой жидкости в плоском слое 

## Марина В. Ефимова

В настоящей статъе рассматривается нелинейная модель конвекиии в приближении ОбербекаБуссинеска, описываюшая плоское совместное движение бинарной смеси и вязкой теплопроводной жндкости с общей поверхностью раздела. Важсно, что продольный градиент температурь и концентрации имеет квадратичную зависимость от координат х. Построено стационарное решение системы.

Ключевые слова: уравнения Обербека-Буссинеска, конвективное движнение, бинарная смесь, установившееся течение.


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